

Randall-Sundrum with Kalb-Ramond field: return of the hierarchy problem?

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CANADA

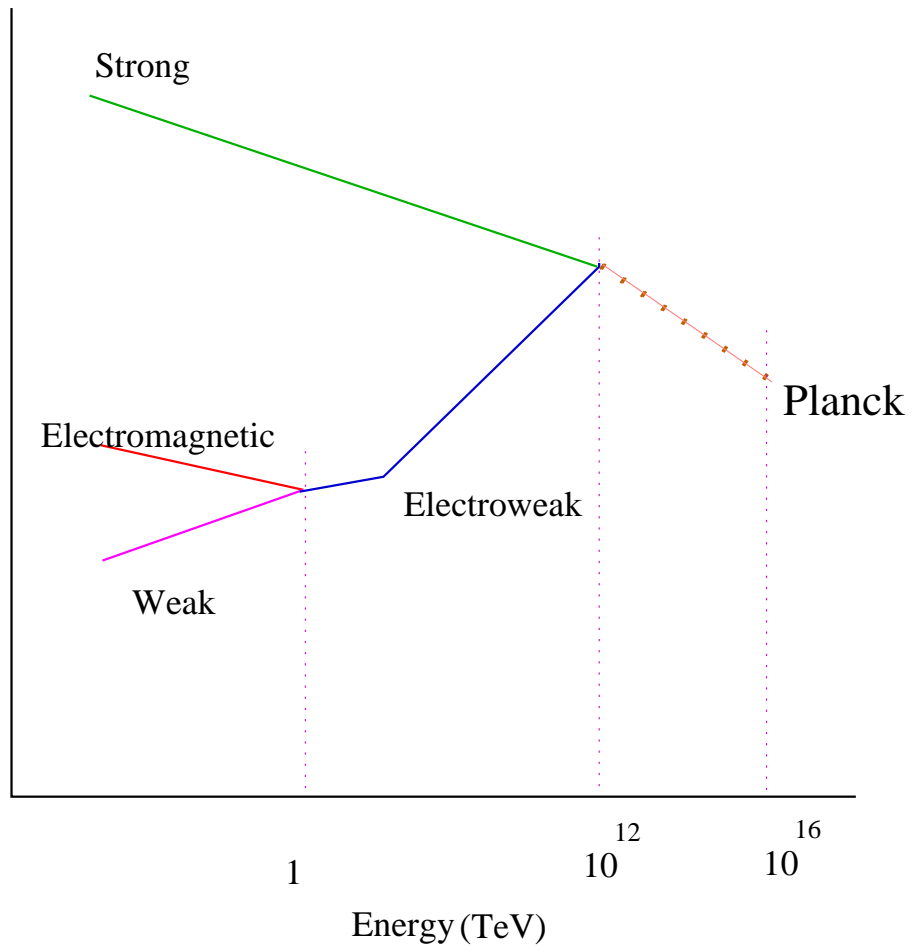
PLAN:

- The hierarchy problem and RS brane world scenario
- The Kalb-Ramond field and the return of the hierarchy problem
- Induced negative 4 dimensional cosmological constant

Reference:

S. Das, A. Dey, S. SenGupta, hep-th/0511257, to appear in Class. Quant. Grav. (Letters)

The hierarchy problem



16 orders of magnitude

TeV Planck scale from Large Extra Dimensions

No of Spacetime dimensions $d = \underbrace{4}_{\text{Observed Standard Model}} + \underbrace{(d-4)}_{\text{Unobserved Gravity}}$
& Gravity

RS Scenario:

Einstein action in d dimensions:

$$S = \frac{c^3}{16\pi G_d} \int d^d x \sqrt{-g_d} \mathcal{R}_d$$

Assume: $ds_d^2 = e^{-A(y)} \underbrace{ds_4^2}_{\text{Observed}} - \underbrace{dy_I dy^I}_{\text{Unobserved}}$
 \uparrow
 non-trivial role of y_I s:

Then:

$$S = \frac{[c^3 \int d^{d-4} y \sqrt{g(y)} e^{-A}]}{16\pi G_d} \int d^4 x \sqrt{-g_d} \mathcal{R}_d = \frac{c^3}{16\pi G_4} \int d^4 x \sqrt{-g_4} \mathcal{R}_4$$

$$G_4 = G_d \left[\int d^{d-4} y \sqrt{g(y)} e^{-A} \right]^{-1}$$

$$M_{Pl(d)}^{d-2} = \left(\frac{\hbar}{c}\right)^{d-4} \frac{M_{Pl(4)}^2}{\int d^{d-4} y \sqrt{g(y)} e^{-A}} \approx \left(\frac{\hbar}{c}\right)^{d-4} k^{d-4} M_{Pl(4)}^2$$

[$A = kL$]

However:

If our universe ('brane') is located at $y = y_0$, metric has conformal factor $\Omega^2 = e^{-A(y_0)}$

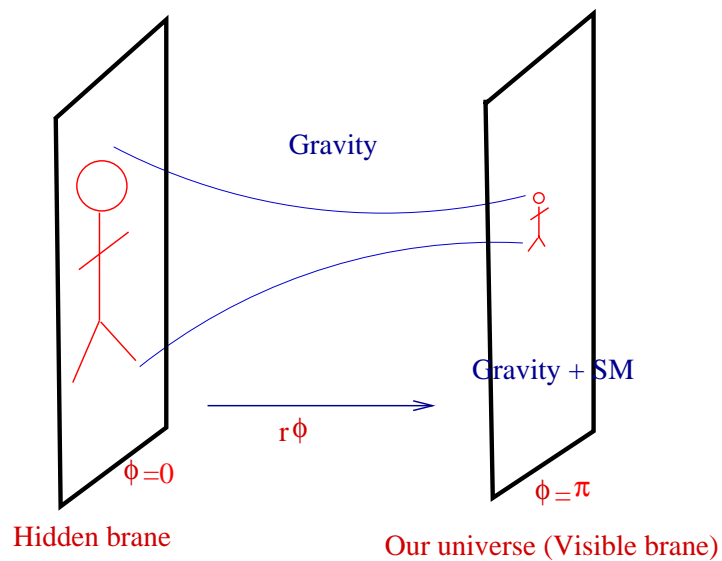
Physical masses $M_{(4)}^{phys} = e^{-A(y_0)} M_{(4)}^{param}$

Thus if $M_{(4)}^{param} c^2 = 10^{16} TeV$ and $A \approx 12$

Then:

$M_{(4)}^{phys} c^2 = 1 TeV$

A Small conformal factor is sufficient to solve hierarchy



Computing the warp factor

Ansatz:

$$ds^2 = e^{-A} \eta_{\mu\nu} dx^\mu dx^\nu - r^2 d\phi^2 \leftarrow \text{extra dim}$$

Action ($M_{Pl(5)} \equiv M, \mathcal{R}_5 = R$):

$$S = S_{Gravity} + S_{vis} + S_{hid}$$

where, $S_{Gravity} = \int d^4x r d\phi \sqrt{-G} [2M^3 R + \underbrace{\Lambda}_{5-d}]$

$$S_{vis} = \int d^4x \sqrt{-g_{vis}} [L_{vis} - V_{vis}]$$

$$S_{hid} = \int d^4x \sqrt{-g_{hid}} [L_{hid} - V_{hid}]$$

Eq. of motion:

$$\frac{3}{2} A'^2 = -\frac{\Lambda}{4M^3} r^2 \quad \left[' = \frac{d}{d\phi} \right]$$

Solution:

$$A = 2kr\phi$$

$$V_{hid} = -V_{vis} = 24M^3 k \quad \left[k = \frac{-\Lambda}{24M^3} \right]$$

Warping

$$\left(\frac{m_H}{m_0}\right)^2 = e^{-2A}|_{\phi=\pi} = e^{-2kr\pi} \approx (10^{-16})^2$$

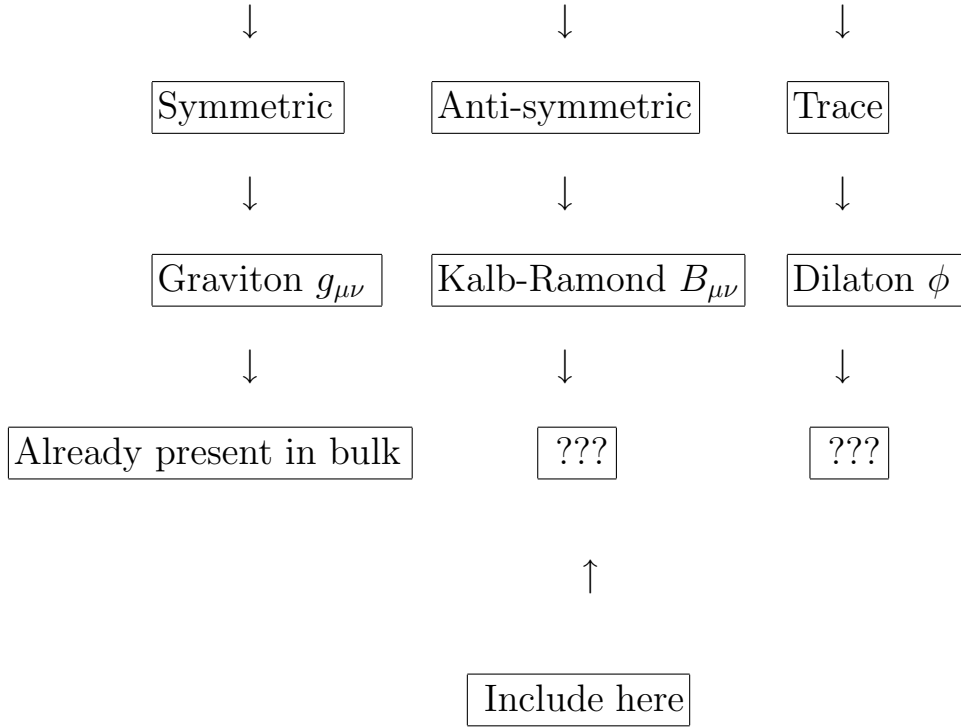
$$\Rightarrow kr = \frac{16}{\pi} \ln(10) = 11.6279\dots \quad \leftarrow \text{RS value}$$

What if there are other fields in the bulk?

Branes in String Theory ? \Rightarrow *D*-branes?

Massless modes in string theory

$$\alpha_{-1}^{\mu} \tilde{\alpha}_{-1}^{\nu} |0; k\rangle \leftarrow \text{Massless}$$



Note:

$$H_{MNL} = \partial_{[M} B_{NL]}$$

Computing the warp factor with bulk KR field

Ansatz:

$$ds^2 = e^{-A} \eta_{\mu\nu} dx^\mu dx^\nu - r^2 d\phi^2 \leftarrow \text{extra dim}$$

Action:

$$S = S_{Gravity} + S_{vis} + S_{hid} + S_{KR}$$

where, $S_{Gravity} = \int d^4x r d\phi \sqrt{-G} [2M^3 R + \underbrace{\Lambda}_{5-d}]$

$$S_{vis} = \int d^4x \sqrt{-g_{vis}} [L_{vis} - V_{vis}]$$

$$S_{hid} = \int d^4x \sqrt{-g_{hid}} [L_{hid} - V_{hid}]$$

$$S_{KR} = \int d^4x r d\phi \sqrt{-G} [-2H_{MNL}H^{MNL}] \leftarrow \text{Additional}$$

Eq. of motion:

$$\frac{3}{2}A'^2 = -\frac{\Lambda}{4M^3} r^2 - \frac{3}{2M^3} g^{\nu\beta} g^{\lambda\gamma} H_{\phi\nu\lambda} H_{\phi\beta\gamma} \leftarrow \text{New}$$

$$\frac{3}{2}(A'^2 - A'') = -\frac{\Lambda}{4M^3} r^2 + \frac{\exp(-2A)}{2M^3} \eta^{\lambda\gamma} [-12\eta^{00} H_{\phi 0\lambda} H_{\phi 0\gamma} + 3\eta^{\nu\beta} H_{\phi\nu\lambda} H_{\phi\beta\gamma}]$$

$$\frac{3}{2}(A'^2 - A'') = -\frac{\Lambda}{4M^3} r^2 + \frac{\exp(-2A)}{2M^3} \eta^{\lambda\gamma} [-12\eta^{ii} H_{\phi i\lambda} H_{\phi i\gamma} + 3\eta^{\nu\beta} H_{\phi\nu\lambda} H_{\phi\beta\gamma}]$$

↑
New

↑
New

Solution:

$$e^{-A} = \frac{\sqrt{b}}{2kr} \cosh(2kr\phi + 2krc)$$

$b \sim$ Energy density of KR field

$$\frac{2kr}{\sqrt{b}} = \cosh(2krc) \quad , \quad \text{such that } A(0) = 1$$

$$c = -\frac{1}{2kr} \tanh^{-1}\left(\frac{V_{hid}}{24M^3k}\right) = -\pi + \frac{1}{2kr} \tanh^{-1}\left(\frac{V_{vis}}{24M^3k}\right)$$

RS Limit:

$$b \rightarrow 0 \quad , \quad c \rightarrow -\infty$$

Warping

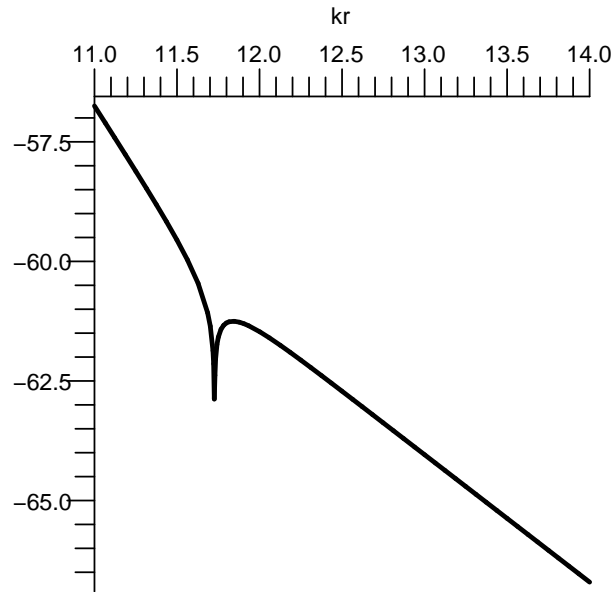
$$\begin{aligned} \left(\frac{m_H}{m_0}\right)^2 &= e^{-2A}|_{\phi=\pi} = \frac{\sqrt{b}}{2kr} \cosh \left[2kr\pi + \cosh^{-1} \frac{2kr}{\sqrt{b}} \right] \leftarrow \text{New warp factor} \\ &= \left[\cosh(2kr\pi) - \sinh(2kr\pi) \sqrt{1 - \frac{b}{(2kr)^2}} \right] \\ &\approx (10^{-16})^2 \end{aligned}$$

Invert to get $b = b(kr)$:

$$b = (2kr)^2 \left[1 - \left(\coth(2kr\pi) - (m_H/m_0)^2 \operatorname{cosech}(2kr\pi) \right)^2 \right]$$

$$\begin{aligned} kr = \text{RS value} &\Rightarrow b = 0 \\ kr > \text{RS value} &\Rightarrow b > 0 \\ kr < \text{RS value} &\Rightarrow b < 0 \leftarrow \text{Unphysical, since then } \left(\frac{m_H}{m_0}\right)^2 \in \mathcal{C} \end{aligned}$$

Therefore, $b > 0$, but how large ?



$\log |b|$ vs kr , for $\frac{m_H}{m_0} = 10^{-16}$

kink = RS value of kr ($b = 0$)

Left of kink: $kr <$ RS value, $b < 0$ (unphysical)

Right of kink: $kr >$ RS value, $b > 0$ (physical)

Any value of $kr >$ RS value solves hierarchy

But

$$b_{max} = 10^{-62}$$

↑

Extreme fine tuning of b !!

Induced 4-dim cosmological constant (0 for RS)

$${}^{(5)}G_{AB} = -\Lambda {}^{(5)}g_{AB} + 8\pi G_5 T_{AB}$$

↓

$$G_{\mu\nu} = -\lambda g_{\mu\nu} + \dots$$

$$\lambda \equiv \frac{1}{2} (kV_{vis} + \Lambda) = -12M^3k [\tanh(2krc) + 1]$$

$$\approx -24M^3k \frac{b}{(4kr)^2} \approx -10^{-63} \neq +10^{-123} \leftarrow \text{accepted value}$$

SUMMARY AND OUTLOOK:

- RS model solves hierarchy problem with gravity and KR field in the bulk
- *BUT* Extreme fine tuning of KR field: 1 part in 10^{62} ← *Return of the hierarchy/fine-tuning problem*
- Wrong sign of induced 4-dim cosmological constant
- Include dilaton, other fields?
- Is RS brane world the answer??