

① "Non-Equilibrium Thermodynamics of Spacetime"

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- What are the states that bh entropy counts?
- ~~Bh~~ Bh entropy or horizon entropy?
- Newton's constant and bh entropy: renormalization?
- ~~is gravity thermodynamical, different?~~

10 yrs old derivation of EE as eq'n of state based on some hypotheses:

① all local causal horizons have a universal entropy density \propto
 $S = \alpha A$.

② Clausius rel'n:

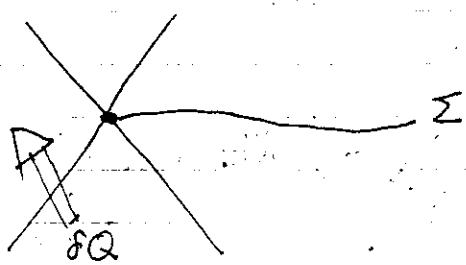
$$\downarrow S = \frac{\delta Q}{T}$$

δQ : boost energy flux across horizon.

$T = \frac{\hbar}{2\pi}$: Unruh temperature

$$\Rightarrow \textcircled{R} R_{ab} - \frac{1}{2} R g_{ab} - \Lambda g_{ab} = 8\pi G T_{ab}$$

Λ : arbitrary, $G = (4\pi\alpha)^{-1}$



Focusing of light rays \rightarrow gravity

Explain it here plus recent lecture one w/ Chris Elving and Raf Guedens incorporate non-equilibrium viscous terms. - incorporate higher curvature corrections to

on horizon: Minkowski spacetime, boost symmetry entropy (what led us to noneqil).

$k^a = (\frac{\partial}{\partial \lambda})^a$; if $\lambda = \text{aff parameter out}$, $\chi = x \partial_t + t \partial_x$; bifurcation plane $x=t=0$.
 $\Rightarrow \chi^a = -\lambda k^a$ $U = \text{null parameter}$

$\chi^a = (\frac{\partial}{\partial \lambda})^a$
 $\lambda = -e^{-u}$
 $\frac{d\lambda}{d\lambda} = -\lambda$
 $H_\chi = \int d\Sigma^a T_{ab} \chi^b$; $d\Sigma^a = d^2 A k^a d\lambda$

$$T_{\Sigma} |0\rangle\langle 0| = z^{-1} \exp(-\frac{2\pi}{\hbar} H_\chi)$$

Unruh effect: the root of bh thermo?

(2)

Curved spacetime: "local boost symmetry"

Killing's eqⁿ: $X_{a;b} + X_{b;a} = 0$

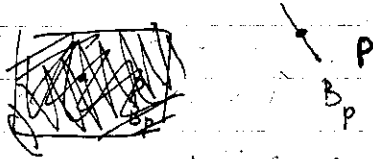
Solⁿ determined by X_a & $X_{a;b}$ at one point.

Minkowski boost: $X_a|_p = 0$ $X_{a;b}|_p =$

t	x	y	z
0	+1	0	0
x	-1	0	0
y	0	0	0
z	0	0	0

$X_{a;b} \perp$ bifurcation plane.

$X_{a;b} X^{a;b} = -2$

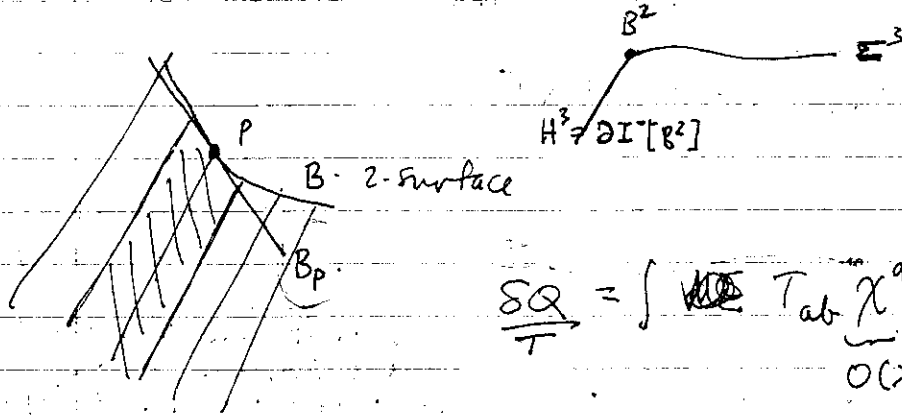


two surface through p

try to solve Killing's eqⁿ in Riemann normal coord's -

→ 2nd order terms vanish

3rd " " cannot be satisfied



$$\frac{\delta Q}{T} = \int_{\Sigma^3} T_{ab} \chi^a d\Sigma^b$$

$O(\lambda)$

$$= \int_{\Sigma^3} T_{ab} k^a k^b (-\lambda) d\lambda d^2A$$

$$\delta A = \int \theta d^2A, \quad \theta = \frac{d \ln(d^2A)}{d\lambda}$$

$$\frac{d\theta}{d\lambda} = -\frac{1}{2}\theta^2 - \sigma_{ab}\sigma^{ab} - R_{ab}k^a k^b$$

3

$\delta_p = 0$ $O(\lambda)$ part ^{must} vanish.

$$\delta S = \alpha \int (R_{ab} k^a k^b + \sigma_{ab} \sigma^{ab}) (-\lambda) d\lambda d^2A$$

$$\delta S = \frac{\delta Q}{T} \Rightarrow R_{ab} k^a k^b + \sigma_{ab} \sigma^{ab} = \frac{2\pi}{h\alpha} T_{ab} k^a k^b \quad \forall k$$

$$\Rightarrow \sigma_{ab} = 0 \quad \& \quad R_{ab} + \Phi g_{ab} = \frac{2\pi}{h\alpha} T_{ab}$$

previously just assumed as "equilibrium", but not so. in any case it's a form of equilibrium... but how quickly?

$$\left. \begin{array}{l} \text{Bianchi} \\ \& \\ T_{ab};^b = 0 \end{array} \right\} \Rightarrow \Phi = -\frac{1}{2} R \bar{\Lambda}$$

$$\frac{2\pi}{h\alpha} = 8\pi G \Rightarrow \boxed{G = \frac{1}{4h\alpha}}$$

Not as quickly σ_{ab} : shear wrt affine parameters

$\hat{\sigma}_{ab}$: s " Killing parameter

$$= \frac{dt}{dU} \sigma_{ab}$$

$$= e^{-U} \sigma_{ab}$$

internal

entropy production: $dS = \frac{\delta Q}{T} + diS$

$$\frac{diS}{dU d^2A} = \frac{1}{T} \cdot 2\zeta \hat{\sigma}_{ab} \hat{\sigma}^{ab}$$

shear viscosity

like fluid exactly.

$$\text{so } \frac{1}{T} \cdot 2\zeta = \alpha$$

$$\Rightarrow \sigma_{ab} \sigma^{ab} (-\lambda) d\lambda = \hat{\sigma}_{ab} \hat{\sigma}^{ab} dU \quad \zeta = \frac{\alpha T}{2} = \frac{h\alpha}{4\pi} = \frac{1}{16\pi G}$$

4

include bulk viscosity?? θ_p^2 : NO. Why not?

~~equilibrium?~~

include curvature correction to entropy.

$$S = \alpha f(R) = \alpha (1 + O(R))$$

$$\delta S = \alpha \int \frac{d}{d\lambda} (f d^2 A) d\lambda$$

$$= \alpha \int (\theta f + \dot{f}) d\lambda d^2 A$$

$$\delta S = \frac{\delta S}{T} \Rightarrow (\theta f + \dot{f})|_p = 0$$

$$\Rightarrow \theta = -\frac{\dot{f}}{f} \neq 0! \quad (Raf)$$

Go thru similar argument & find

$$f R_{ab} - f_{;ab} + \frac{3}{2} f^{-1} f_{,a} f_{,b} + \Psi g_{ab} = \frac{2\pi}{4\pi} T_{ab}$$

$$\Psi = \square f - \frac{1}{2} \mathcal{L} - \Theta, \quad \frac{d\mathcal{L}}{dR} = f \text{ etc.}$$

$$\left. \begin{array}{l} \text{Bianchi} \\ * \\ T_{ab;b} = 0 \end{array} \right\} \Rightarrow \left(\frac{3}{2} f^{-1} f_{,a} f_{,b} \right)_{;a} = \Theta_{,b} \text{ some } \Theta$$

NOT TRUE!

5

Bulk viscosity term

$$\frac{d_i S}{d\omega d^2 A} = \frac{3}{2} \alpha f \hat{\theta}^2$$

$$\begin{aligned} \Rightarrow \text{Bulk viscosity } \eta &= \frac{3}{2} \alpha f T \\ &= \frac{3}{16\pi G} f. \end{aligned}$$

Then derive Einstein eq'n ~~in presence~~ associated w/ Lagrangian, that yields $S = \alpha f A \left(\propto \frac{d\mathcal{L}}{dR} \right)$.

$$\frac{d\mathcal{L}}{dR} \quad (\text{Wald entropy})$$

Comments:

① Clausius + 1st law \rightarrow EDS.
 \downarrow
 $\nabla_a T^{ab} = 0$.

② viscosity a phenomenological coefficient?

Is, in terms of α . e.g. shear $\eta = \frac{t_{\alpha}}{4\pi}$.

[Derivation in string theory via AdS/CFT?]

③ Curvature corrections meaningful? If imagine minimum size ϵ of region for argument, killing approximation fails ... ok if coeff of curvature term large enough ($> \epsilon^2$).

6

④ Other types of curvatures terms, eg. Lovelock?

The End